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Abstract: Gamma-ray bursts (GRBs) are the most luminous objects known: they outshine their host galaxies, making them ideal candidates for probing large-scale structure. Our aim is to determine the Spatial Two-Point Correlation Function of the GRBs with spectroscopic redshifts. We used all observations till 4 June 2021, and we also split the data according to the origin of the redshift (afterglow or host galaxy). We analyzed the possible correlation between the GRB redshift and sky position. There is only one region, called the Faraway GRB Patch, where nine distant GRBs shows deviation from the randomness with an $\approx 1\%$ significance, showing that the sky and the radial component of the GRB distribution could be factorized and both can be determined independently. The GRBs' cumulative monthly distribution was used to show that it is impossible to synthesize the Sky Exposure Function, even from the perfect observational logs. We estimated the Sky Exposure Function using the Gaussian kernel, and with the radial distribution, it was used for the estimation of the Spatial Two-Point Correlation Function. Monte Carlo simulations were used to determine the Poissonian errors. Our result shows that the current Spatial Two-Point Correlation Function estimations are consistent with zero for the full and the afterglow/host galaxy datasets.

Keywords: gamma-ray burst: general; gamma-rays: general; methods: data analysis; methods: statistical; cosmology: large-scale structure of the Universe; cosmology: observations

1. Introduction

Gamma-ray bursts (GRBs) outshine their host galaxies, making them ideal candidates for probing large-scale structure, as they can be detected up to very high redshifts [1–6]. According to the Cosmological Principle, the Universe is spatially homogeneous and isotropic on a large scale: assuming that GRBs follow the distribution of baryonic matter, we would expect homogeneity and isotropy for GRBs as well. Hence, the spatial distribution of the GRBs allows us to test matter's distribution in the Universe.

The paper is organized as follows: Section 2 describes our dataset; Section 3 overviews the previous statistical analyses and results regarding the GRBs' distribution. In Section 4, we determine the Sky Exposure Function for the GRBs with spectroscopic redshift. In Section 5, the Spatial Two-Point Correlation Function is approximated, while Section 6 discusses the results, and Section 7 summarizes.

2. Data Selection

The sample of GRBs with measured position, optical afterglow, and spectroscopic redshift is created largely of bursts detected by NASA's Swift and Fermi experiment. Here,



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we used the same dataset of Horvath et al. [7], which used redshifts from the GRBOX database (Gamma-Ray Burst Online Index) database (the datasets were derived from different sources in the public domain: https://www.astro.caltech.edu/grbox/grbox.php, (accessed on 4 June 2021)). The GRBOX database is heavily based on the notes of the Gamma-ray Coordination Network (http://gcn.gsfc.nasa.gov/, (accessed on 4 June 2021)) and on the public dataset of Joachim Greiner (http://www.mpe.mpg.de/~jcg/grbgen.html, (accessed on 4 June 2021)), which also contains much information about almost every GRB observed by any instrument.

The GRBOX catalog database (published by the Caltech Astronomy Department) has been carefully reviewed by its creators to avoid observations with large systematic uncertainties. The redshift data in Horvath et al. [7] contain spectroscopic redshifts only. Photometric redshifts, as well as redshift estimations based on Ly- α limits are excluded as they introduce large (typically more than several hundreds of Mpc) radial distance uncertainties. By November 2013, the redshifts of 361 GRBs had been determined, and as of March 2018, the redshifts of 487 GRBs have been measured. Here, we expand the data of Horvath et al. [7] with new observations, till 4 June 2021. Greiner's table and the relevant Gamma-ray Coordination Network messages were also used to enhance the data. This process resulted in a total of 522 GRBs with precise position and spectroscopic redshift. The redshift's optical observation type according to afterglow or host galaxy detection is registered as well. The galactic distribution of these 522 GRBs is shown on Figure 1, according to the redshift origin (afterglow or host galaxy measurement).

This GRB positional set is the largest available, worthwhile for a large-scale homogeneity and isotropy check. One should remark that after 2006, the well-localized GRBs rate was \approx 120/year, but later, it decreased almost with an exponential decay: every year, the results were only \approx 90% of the previous year's redshift observations [7]. This trend was halted only in recent years by a few optical campaigns with really committed observers. However, the current rate of the new redshift observations is low, which makes the chance of the substantial statistical improvements low in the near future.

We calculated and used the comoving distances based on the observed spectroscopic redshifts. Following the results of Planck Collaboration et al. [8], we took $\Omega_m = 0.3089$, $\Omega_{\Lambda} = 0.686$, and H = 67.74 km/s for the calculations. As we are concentrating on the large-scale distribution of the GRBs, we omitted the nearest 500 Mpc observations, this length being $\approx 2 \times$ the BAO distance. In Figure 2, the comoving distance distribution is shown for our GRBs.



Figure 1. Galactic distribution of 522 GRBs with measured spectroscopic redshift till 4 June 2021. The disk of the Galaxy hinders the optical follow-up. The blue/red colors mark the afterglow/host galaxy origin of the redshift.

It should be remarked that the observational errors in the redshift measurements are small, usually in the $\delta z \approx 10^{-3}$ range. We used the errors when present in the observation record and estimated with the last digit of the *z* value for the rest. The radial errors follow an approximately log-normal distribution, with 0.31, 1.46, and 3.16 Mpc for the first quartile, median, and third quartile. The distribution does not depend on the afterglow/hosts distinction. For the full-3D positional errors, the reported positional uncertainties were also considered, producing slightly larger 0.35, 1.65, and 3.60 Mpc values for the first quartile, median, and third quartile. These errors do not affect our distance calculations practically.



Figure 2. The comoving radial distance distribution of all GRBs in our dataset (except the nearest 500 Mpc observations). There is almost a linear trend between 4 and 6 Gpc ($z \approx 1.3-2.5$), where the GRB density is the highest. The radial distribution of the Faraway GRB Patch (q.v. Section 4.1) is also shown: here, all distances are from afterglow observations only.

3. Large-Scale Anomalies Observed Earlier

3.1. Compton Gamma-Ray Observatory BATSE's Observations

The first indirect observational proof for the GRBs' cosmological origin was reported by Meegan et al. [9]: after the first 2 years of BATSE observations, the observed GRBs showed an almost isotropic distribution without any concentration around the Galactic plane.

Briggs [10] proposed the Rayleigh–Watson and the Bingham statistics to test the point source distribution's dipole and quadrupole anisotropies. Briggs et al. [11] used dipole and quadrupole statistics for the BATSE GRBs' sky distribution and found no deviation from the large-scale isotropy at the 0.9σ level for the dipole and the 0.3σ level for the quadrupole term. These results strengthened the cosmological origin model of the sources.

Using angular positional measurements only from the BATSE experiment, Balázs et al. [12] found that the sky distributions of short and long GRBs are different. They compared the quadrupole term, which is non-zero, with a probability of 99.97%.

The very short (with $T_{90} < 0.1$ s) GRBs' sky distribution have been shown to be anisotropic by Cline et al. [13] and Cline [14] with a p = 0.00016 level (the non-uniformity in the sky exposure was neglected).

Considering the non-uniform BATSE sky exposure, Chen and Hakkila [15] analyzed the sky exposure regarding the Two-Point Angular Correlation Function and found that the sky exposure effects are small. However, the non-uniform sky exposure complicated all detailed statistics, so it should have been taken care of.

A BATSE-modified count-in-cell method, which compensates for the well-known Sky Exposure Function, was applied on the GRB data: it was found that the distribution of the intermediate class (where $2 \text{ s} < T_{90} < 10 \text{ s}$) is not isotropic on the 99.3% confidence level [16]. Mészáros et al. [17] used the method of spherical harmonics for the investigation of the angular distribution of the same intermediate group and found that it shows an intrinsic anisotropy at the 97% significance level. Litvin et al. [18] confirmed the anisotropy for the short bursts with a higher 99.99% significance level, while for the intermediate class, anisotropy only, a lower 99.89% significance level was obtained.

On the 2–5° angular scale, Magliocchetti et al. [19] found significant correlation among the short GRBs at a level of $\sim 2\sigma$. Tanvir et al. [20] confirmed a correlation between the short GRBs' position and the positions of galaxies in the local Universe at a 99.9% level.

Using the full BATSE data, it has also been shown using extensive combined statistics (using Voronoi tessellation, minimal spanning tree, and multifractal spectrum) that the short GRBs deviate significantly from the full randomness at a 99.90–99.98% level, while the intermediate group showed a weaker, but still significant deviation from the full isotropy at a 98.51% level [21]. Extensive Monte Carlo simulations were applied to compensate for the BATSE's sky exposure function.

3.2. The Hercules-Corona Borealis Great Wall

The largest observed structure in the Universe so far was identified as a large GRB cluster at redshift $z \approx 2$ in the general direction of the constellations of Hercules and Corona Borealis [22].

The size of the Hercules-Corona Borealis Great Wall is estimated to be about 2–3 Gpc across (Figure 3); the uncertainties are mainly due to the sparse sampling. This clustering's scale is quite large: if it reflects the underlying distribution of matter, the cluster is big enough to question standard assumptions about Universal homogeneity and isotropy.



Figure 3. The distribution of GRBs in the redshift range of $1.6 < z \le 2.1$, using current data. The Hercules-Corona Borealis Great Wall's GRBs [7] are highlighted in violet.

Horvath et al. [22] used 283 GRB with redshifts to investigate the spatial distribution of GRBs. They assumed that the sky exposure is independent of *z*; therefore, they subdivided the GRB sample by redshift. Angular tests were applied on nine groups based on redshift/distance, using *k*-th nearest neighbor analysis and the bootstrap point radius method on their dataset. Nearest neighbor tests identified pairing consistent with the previously found large, loose GRB cluster in the redshift range $1.6 < z \leq 2.1$.

The observed angular excess cannot be entirely attributed to the known selection biases, making its existence due to chance unlikely. The significance of this cluster can be measured using the bootstrap point–radius method where one selects random locations on the celestial sphere and finds how many of the points lie within a circle of a predefined angular radius. The analysis can be performed with the clustered 44 GRBs and also with 44 random samples from the full observed dataset [23]. The probability that this clustering is random is $p = 1.6 \times 10^{-4}$. The direction of the cluster is $\approx l = 88^{\circ}$, $b = 63^{\circ}$.

The initial discovery was supported by subsequent observations [7,23], although the exact nature of the structure is still unknown partially due to the sparse data. Horvath et al. [7] divided the data into six radial GRB groups and found the 1.6 < z < 2.1 group to be the most anisotropic. They suggested incorporating the sky exposure for a more detailed spatial analysis—this is what we used for the Spatial Two-Point Correlation Function calculations in Section 4. In Figure 3, the relevant $1.6 < z \leq 2.1$ range is shown in our current GRB data.

3.3. The Giant GRB Ring

Balázs et al. [24], motivated by the Hercules-Corona Borealis Great Wall, analyzed the *k*-th nearest neighbor in the sample further. Instead of the slices in the redshift range, the *k*-th Next Neighbor Statistics was used to determine the spatial density of the GRBs. Multiple *k* values were used to calculate the mean and variance of the local density at every GRB location and, it was found that in the cases of k = 8, 10, 12, 14, the grouping of nine objects deviates significantly from the random case.

During the analysis, a large regular formation of GRBs was found: the Giant GRB Ring is displayed by nine GRBs, with an angular major/minor diameter of $43^{\circ}/30^{\circ}$, at a comoving distance of ≈ 2.77 Gpc in the 0.78 < z < 0.86 redshift range, with a probability of 2×10^{-6} of being the result of a random fluctuation only [25]. The size is ≈ 1.7 Gpc, a bit smaller than the Hercules-Corona Borealis Great Wall, but larger than the large quasar group ($1240 \times 640 \times 370$ Mpc, [26]). In Figure 4, the Giant GRB Ring's 0.78 < z < 0.86 redshift range is shown using current GRB data.

Regarding the Giant GRB Ring, Balázs et al. [25] used different statistical methods to analyze the statistical properties of the spatial point process displayed by GRBs of known redshift. They used kernel-smoothing for the the joint probability density estimation, while the independence of the angular and radial variables of the GRBs was supposed. With 1502 random datasets, they found only three ring-like patterns having similar parameters as the original one appearing in the Ring; thus, the Giant GRB Ring's existence is a lowprobability random event. They also warned that the large-scale GRB distribution does not necessarily reflect the large-scale distribution of the matter.



Figure 4. The Giant GRB Ring's 0.78 < z < 0.86 redshift range in the current data. The ring-like structure of objects is in the lower right part of the frames, highlighted in violet [26].

3.4. Large-Scale Distribution and Physical Properties of GRBs

There are at least two different types of GRB central engines: the merging of compact objects, such as neutron stars and/or black holes, will probably produce short GRBs, confirmed by, e.g., the kilonova observation of Tanvir et al. [27] and the GRB170817A event [28]. Collapsing massive stellar objects are connected to the long GRBs [29]. There are some indications for the existence of a third, intermediate group [30–34], and using different observed parameters and spectral models, the classification could be extended [35].

The early CGRO BATSE observation anisotropy investigations [12,13,16–19,21] have shown that the major GRB groups, especially the short and the long one, show different sky distribution. They are important from the sky exposure point of view: if events with different observed length have the same observation probability, the anisotropy effects should be real. Therefore, any correlation between the physical parameters and the sky distribution is important.

Tarnopolski [36] used the sky distribution of 1669 Fermi/GBM GRBs for the binomial test, nearest neighbor, fractal dimension, dipole and quadrupole analysis, and the two-point angular correlation tests. The results showed that short GRBs are distributed anisotropically in the sky with a probability of 99.98%, while the long GRBs showed no anisotropy. In Bagoly et al. [37], 361 GRBs with position and redshift were analyzed, and the large-scale homogeneous and isotropic distribution was checked using nearest neighbor tests and the two-point correlation function. Shirokov et al. [38] analyzed the spatial distribution of 384 Swift GRBs with known redshifts using conditional density and pairwise distance methods. They estimated a fractal dimensionality of D = 2.55 on scales of 2–6 Gpc.

Andrade et al. [39] tested the sky distribution isotropy of the FERMI-GBM catalogue's 2626 GRBs. The large positional uncertainties in the positions prevented the applied twopoint angular correlation function from detecting any statistical anisotropy, both for the long and short GRBs.

Wang et al. [40] analyzed some intrinsic properties of 6289 GRBs using prompt and afterglow parameters, searching for correlations between the parameters and the classification of the GRBs.

Cao et al. [41] used the Platinum GRB data compilation to create standardizable events in the Dainotti correlation space. The obtained GRB cosmological parameter limits are consistent with the baryon acoustic oscillation (BAO) and Hubble parameter [H(z)] data. In Cao et al. [42], Dainotti-correlated gamma-ray burst were standardized and used to constrain cosmological model parameters. The results provide limits on the non-relativistic matter density, slightly favoring non-zero spatial curvature.

Řípa and Shafieloo [43] used the Fermi/GBM GRB catalog to test the correlation between GRBs' durations, fluences, peak fluxes measured in different energy, and the sky positions of the GRBs. Řípa and Shafieloo [44] expanded the data with the BATSE and Swift BAT GRBs. The tests showed no correlation between the sky positions and the physical properties of the GRBs.

Recently, Horvath et al. [45] investigated the connection between the GRBs' durations and redshifts/radial distances and also found no correlation between them.

These results are consistent with a view that the GRBs physical parameters are independent of their position in the Universe.

3.5. Quasar Positions and Polarizations

Although we do not know precisely how the GRBs and the visible matter distributions are related, the quasars' distribution in the Universe could give some hints about the baryonic matter distribution. The quasars' high luminosities and early emission make them appropriate for large-scale surveys besides the GRBs.

Jain et al. [46] analyzed the polarization of quasar emission with parallel transport techniques. The statistical results support the existence of the large-scale alignment effect in the z > 1 quasar optical polarization observations. Pelgrims and Hutsemékers [47] found significant large-scale alignments of polarization vectors among quasars using 2D and 3D

statistical methods. Using 71 quasars at redshifts 1.0 < z < 1.8, Friday et al. [48] reported correlated orientations of their optical and radio polarization on Gpc scales.

Clowes et al. [26] analyzed the spatial distribution of quasars, claiming the existence of a structure with a size above 1 Gpc. Marchã and Browne [49] used two-point correlation functions to analyze the clustering of Fermi-selected BL Lac objects and Fermi-selected flat spectrum radio quasars. Separating the BL Lac-like objects with machine learning classifications based only on gamma-ray information, they found evidence for large-scale clustering amongst them. As an explanation for the clustering, they proposed that there are indeed large volumes of space in which AGN axes are aligned.

Lopez et al. [50] analyzed the distribution of the intervening MgII absorbers in the spectra of background quasars and found a "Giant Arc on the Sky" at $z \sim 0.8$. Its size is \sim 1 Gpc and appears to be almost symmetrical on the sky.

3.6. Theoretical Remarks about the Anisotropies

The existence of the above-mentioned anomalies raises the question about their origin. Without full analysis, we list some theories. See [51] for an overview of the standard Λ CDM model's tensions problem and other current cosmological challenges.

One source could be some real cosmic anisotropy, e.g., Ciarcelluti [52] studied the anisotropic cosmic expansions with electrodynamics, predicting the appearance or variation of the polarization of electromagnetic radiation regarding these regions. Migkas et al. [53] analyzed 570 clusters with X-ray, microwave, and infrared observations to create scaling relations. Using all available distance information, they detected an apparent 9% dipole-like spatial variation in the local H_0 at (l, b) = (280, 15) on the sky. The result could also be attributed to a 900 km/s bulk flow. Any variation in the H_0 will create different radial scales and would produce an anisotropy signal.

There could be several astrophysical reasons for the clustering of GRBs [7]. The low metallicity massive stars are thought to be long GRB progenitors as they generally appear in faint, blue, low-mass star-forming galaxies [54,55] and in the bright regions of the host galaxy [56–58]. Short GRBs are not expected to trace star-forming activity due to the long inspiralling time of the binary. This notwithstanding, Palmerio et al. [55] claims that long GRBs are not good tracers of star formation, based on the result of the Swift GRB Host Galaxy Legacy Survey, covering 0.03 < z < 6.3. Bagoly et al. [59] studied the distribution of the starburst galaxies from the Millennium XXL Simulation [60]. With Markov Chain Monte Carlo simulations, the connection between the large-scale structures and the starburst galaxies' groups distribution were checked at the Giant GRB Ring's z = 0.82 distance. Using the same simulation, Balázs et al. [24] demonstrated that the spatial distribution of galaxies with a normal vs. high star-formation rate is probably different. Hence, speculations about some fluctuation or a wave in star formation rate is too simple to describe the full view.

4. The Sky Exposure Function

4.1. Correlation between the Redshift and Sky Position

We know that the observational probability of a GRB differ on the sky, due to the various effects (geometrical, satellite operation, optical follow-ups). As we saw in Section 3.4, no known correlations have been found between the GRBs' direction and their physical properties.

For the reconstruction of the full three-dimensional GRB distribution, the main question is whether the sky distribution is independent of the redshift (comoving radial distance) or not. We tested this correlation in the following way: around each GRB, the radial distance distribution of the neighboring GRBs within an α -sized spherical cap were determined. The Mann-Whitney U-test's Z score was determined between the full and the local dataset using the full radial distribution (i.e., Figure 2). As the full dataset contains the local data, as well, we performed Monte Carlo randomization and determined the Z score distribution for a total random local dataset. The whole process was repeated for the $\alpha = (5^\circ, 50^\circ)$ range for 1000 random datasets each. The result are shown in Figure 5: all the 99%, 99.5% and 99.9% values of the Monte Carlo simulations, as well as the smallest Z score from the real data (the second smallest values are above the 99% curve of the Monte Carlo simulations).

This method selects an interesting region with a spherical cap size of $12-14^{\circ}$, with a corresponding random probability of $\approx 1\%$. This signal is from one region, the Faraway GRB Patch (Figure 6), where nine distant GRBs are concentrated (Figure 2). All redshifts are from afterglow observations for the Faraway GRB Patch.



Figure 5. Mann-Whitney U-test's Z scores of the real data's lowest values (blue) and the 99%, 99.5%, and 99.9% values of the Monte Carlo simulations. The real data's second smallest Z scores are greater than the 99% curve of the Monte Carlo simulations, therefore are not shown for clarity. One should observe the $12-14^{\circ}$ region with a random probability of $\approx 1\%$.



Figure 6. The sky distribution of the Faraway GRB Patch among all GRBs. The angular size is quite small, below the Sky Exposure Function's optimal smoothing Gauss kernel size (Section 4.3).

The \approx 1% random probability of the Faraway GRB Patch makes it interesting; however, it is large enough not to be considered a real strong signal. Another factor is that we repeated the test many times, and there was a correlation between the tests, as the same overlapping data were used. For the 15° spherical cap, there are \approx 60 separate regions on the sky, i.e., this should be the minimum value considered for the correction factor.

Our results show that the sky and the radial component of the GRB distribution could be factorized and both can be determined independently.

4.2. Optical Observational Factors

The detection probability of the redshift-determined GRBs on the sky is a multiparameter event. The detection probability depends on the various trigger and detection conditions of the space instruments, the optical follow-up procedure, etc.

The simplest models for the sky exposure could be created, e.g., using the pointing logs of the satellite. From a theoretical point of view, one can imagine that it is possible to integrate the satellite's FOV timeline (e.g., Swift and/or Fermi GBM pointings) and simulate all the (known) triggering effects, including the background changes. Beside this, one would still need to combine it with a reliable model for selection effects including the impact of galactic extinction and the various optical observational effects and constraints (e.g., timing, sky visibility, telescope/instrument availability, observer interest, etc.). Here, we investigate the available dataset and show that this is practically not feasible as it cannot be done with the required precision due to both technical and human factors.

In Figure 7, we plot the N(>RA) cumulative number of GRBs as a function of the RA right ascension, for every month. All 522 GRBs from the current data were used. Here, we split the data according to the observations: redshift determined from afterglow means that prompt (or almost prompt) observations was performed—there are 354 such GRBs in our sample. Host galaxy observations could be done later; therefore, it is possible that the observational limits and strategies are different. We have 168 redshifts from host galaxy measurements. Campaigns and dedicated projects were (and hopefully will be) organized to measure the redshift of the known host galaxies: sometimes, this is done years after the GRB. The dichotomy between the afterglow and host galaxy determined redshift grouping is boosted further by the fact that host-based observations are rare above $z \approx 3$, i.e., their average redshift is lower.

One would expect a smoothly varying distribution, considering the various influences of the telescope geographic position and availability, weather, projects, etc. The afterglow and host galaxy measurements follow approximately the same pattern with quite similar yearly modulation, except the afterglow's October distribution. The extra \approx 10–15 observations were created probably by other factors, such as a dedicated rapid response afterglow fall campaign.



Figure 7. The *N*(>*RA*) cumulative number of GRBs as a function of the *RA* right ascension, for every month. Current data were used. The data were divided according to the observational method (afterglow or host galaxy). The afterglow and host galaxy measurements follow approximately the same pattern, except the afterglow's October distribution. It was created probably by other factors, such as a dedicated rapid response fall campaign.

4.3. Kernel-Based Estimation of the Sky Exposure Function

As the synthetic way is not quite viable, we can use only the observational data themselves to derive the detection probability on the sky. This is not a paradox: assuming the sky distribution independent of the redshift (see Section 4.1), the dataset-averaged sky detection probability can be approximated and used as a reference.

In Bagoly et al. [61], the optimal empirical Sky Exposure Function was reconstructed with different techniques. The Swift's exposure map was used as a model for the GRB-like angular distribution. The efficiency of the fixed and the adaptive width Gaussian kernel, the Delaunay Tessellation Field Estimator, and the Voronoi Diagram Field Estimator methods were compared on a random field, similar to the Swift's exposure map. It was determined that the fixed-width Gaussian kernel is the optimal among the methods for approximating the exposure map and for determining the optimal width.

In Li and Lin [62], Bagoly et al. [37], as well as Ukwatta et al. [63], kernel-based methods were applied in the reconstruction of the GRBs' empirical Sky Exposure Function. We applied Gaussian kernel-smoothing on the current GRB data to obtain the empirical Sky Exposure Function. In Figure 8, one can observe the optimal kernel-smoothed Sky Exposure Function. There is a factor of \gtrsim 3 between the minimum and maximum, and the galactic disk is visible. The optimal kernel sizes for the typical few hundreds GRBs are in the 20–40° range [61], clearly bigger than the width of the galactic optical extinction range (cf. Figure 1). Thus, kernels will produce a smoother sky exposure than the real one, resulting in a lower statistical power for structures' detection near and below the kernel size.



Figure 8. The optimal Gaussian-kernel-smoothed Sky Exposure Function of the current GRB data. A factor of \gtrsim 3 is between the minimum and maximum, and the galactic disk is also visible (cf. Figure 1 and Figure 4 of Balázs et al. [25]).

5. The Spatial Two-Point Correlation Function

Using the empirical Sky Exposure Function of Figure 8 and the radial distribution (cf. Figure 2) together, we determined the general Spatial Two-point Correlation Function, which is widely used measure to quantify the large-scale structure. We used the Landy–Szalay method [64] as the optimal tool for a known partial vignetting and reduced observational probabilities. For the $\xi(r)$ Spatial Two-Point Correlation Function, one should generate random points according to the spatial detection probability, e.g., here, using the radial distribution and sky exposure. The *DD*, *DR*, and *RR* correlations were calculated between the data–data, data–random, and random–random points. The $\xi(r)$ Spatial Two-Point Correlation Function will be given by $\xi(r) = (DD - 2DR + RR)/RR$.

We used 10,000 random points for the calculations and also used 100 synthetic dataset following the same distribution as the data to create Monte Carlo simulations. The $\xi(r)$ from the random datasets were used to determine the Poissonian errors in the calculations.

In Figure 9, the $\xi(r)$ of the full dataset are shown, with the mean and $\pm 3\sigma$ errors. The 100 synthetic datasets' mean should be 0, and the $\pm 3\sigma$ error lines mark the noise. All GRB ξ values are within the errors, and no correlation structure is observable.

We repeated the same calculations for $\xi(r)$ using only the afterglow or host galaxy distances. In Figures 10 and 11, the corresponding $\xi(r)$ functions are displayed, with the accompanying errors. In both cases, the GRB $\xi(r)$ values are within the 3σ errors.



Figure 9. The Spatial Two-Point Correlation Function of the full dataset (blue) with the MC simulations' mean and $\pm 3\sigma$ errors (red lines).



Figure 10. The Spatial Two-Point Correlation Function of the afterglow-based redshift fraction (blue) with the corresponding MC simulations' mean and $\pm 3\sigma$ errors (red lines).



Figure 11. The Spatial Two-Point Correlation Function of the host galaxy based redshift fraction (blue) with the corresponding MC simulations' mean and $\pm 3\sigma$ errors (red lines).

6. Discussion

Kernels will produce a smoother Sky Exposure Function than the real one, smoothing out the distribution and resulting in a lower statistical power for structures' detection near and below the kernel size. This is possibly why the signals of the Hercules-Corona Borealis Great Wall and/or the Giant GRB Ring are below the 3σ lines: they are buried in the noise. It is not surprising that the spatial two-point correlation is not the optimal tool for such huge scales: besides the low-count Poisson noise, the smoothed Sky Exposure Function will reduce the sensitivity of the $\xi(r)$ function to large structures, i.e., structures will be smoothed out on the sky and the statistics will give lower significance. Kernel-generated sky exposure and the related random catalogs will contain some part of the smoothed structures, making the signal smaller in the Landy–Szalay method.

These effects will distort the distribution quite differently at different comoving distances because of the angular distance scale's saturation, e.g., for z = 0.8, the comoving distance is ≈ 2.7 Gpc and the transverse scale is ≈ 530 Mpc/20°, while z = 1.7 the comoving distance is ≈ 4.5 Gpc with a transverse scale of ≈ 600 Mpc/20°. Hence, the kernel-smoothing is suboptimal as it introduces distortions similar to the redshift space distortions: here, transversal smoothing is experienced, while in the classic redshift measurements, the peculiar velocity components along the line of sight will smear the signal.

One should consider that in Bagoly et al. [61], the 2-year Swift exposure map was used as a benchmark. It is a relatively smooth function due to the large (>90°) BAT aperture, and the best method was selected using this objective function. In Figure 1, one can see that the galactic extinction's effect is acting on a much smaller scale; hence, the best kernel method will definitely oversmooth that region. The low density of the events will further augment the problem.

Clearly, better algorithms should be developed for more precise calculations, such as nearest neighbor analysis-based [36] methods and similar tools, which incorporate the transverse distortion effect mentioned earlier.

We should remark that Bagoly et al. [37] used the Spatial Two-Point Correlation Function to identify the interesting GRB 020819B and GRB 050803 pair. In their data,

both GRBs had a spectroscopic redshift, and they caused a strong peak in Spatial Two-Point Correlation Function around 56 Mpc, with a probability of 0.00996 being a random fluctuation. Quite interestingly, both GRBs' redshift values were questioned later in the literature, more than a decade later after the events. The Optically Unbiased GRB Host (TOUGH) survey [65] measured z = 3.5 photometric redshift for the host galaxy of GRB 050803, unlike the previous preliminary z = 0.422 value, and thus, we omitted this GRB from the dataset. Perley et al. [66] analyzed the purported spiral host galaxy of GRB 020819B at z = 0.41 and used VLT MUSE and X-shooter observations to demonstrate that it is an unrelated foreground galaxy. They identified a possible radio afterglow with a star-forming galaxy at z = 1.9621, which is consistent with the observations. Hence, as both previous estimated distances are probably wrong and the new values are differ significantly, the peak observed in Bagoly et al. [37] obviously should not appear in these results.

7. Summary

We overviewed the previous statistical analyses and results regarding the GRBs' sky and spatial distribution, and in Section 4, we determined some of the various factors influencing the optical observations.

We analyzed the possible correlation between the GRB redshift and sky position. The GRB distance distribution in a spherical cap was compared with the full distribution using the Mann-Whitney U test's Z score. There was only one region, called the Faraway GRB Patch, where nine distant GRBs showed a deviation from the randomness with an $\approx 1\%$ probability. This show that the sky and the radial component of the GRB distribution could be factorized and both can be determined independently.

The GRBs' cumulative monthly N > (RA) functions show an extra component in the afterglow-determined redshifts' October distribution. It was probably originated from some dedicated rapid response fall campaign and shows that it is not possible to create the Sky Exposure Function even from the perfect observational logs.

We used the Gaussian kernel to obtain the estimation of the Sky Exposure Function for our 522 GRB. This distribution was used with the radial distribution for the Spatial Two-Point Correlation Function $\xi(r)$. Monte Carlo simulations were used to determine the Poissonian errors, while the redshift origin (afterglow or host galaxy) factor was applied to split the data. In Figures 9–11, one can see that the current $\xi(r)$ estimations are consistent with zero for the full and the afterglow/host galaxy datasets.

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